## [Contribution from the Cryogenic Laboratory and the Department of Chemistry, The Ohio State University]

## Vapor Pressures of Inorganic Substances. IX. Gallium ${ }^{1}$

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The vapor pressure of gallium has been determined by measuring the rate of effusion of the vapor through an orifice. $\Delta H_{0}^{0}$, the heat of sublimation of gallium at absolute zero, is $65.6 \pm 0.5 \mathrm{kcal}$. The vapor pressure data for gallium are represented by equation 8 in the text. The normal boiling point from this equation is $2500^{\circ} \mathrm{K}$.

## Introduction

The only previous measurements of the vapor pressure of gallium were made by Harteck, ${ }^{2}$ but since the purification of gallium is difficult and since considerable technological improvements in the preparation of gallium have been made in recent times, the vapor pressure of gallium has been redetermined. Furthermore, the possible use of gallium in high temperature heat engines gives added importance to the accumulation of accurate thermal data.

## Experimental Method

Apparatus and Procedure.-The Knudsen method, the determination of the rate at which metal vapor at its equilibrium pressure effuses through an orifice, has been used throughout this research. The quartz vapor cell and the induction heating employed have been previously described. ${ }^{3}$
It was discovered that tantalum reacted with gallium at high temperatures. Consequently a quartz liner was made up for the tantalum Knudsen cell to contain the gallium. A convenient method of loading the Knudsen cell was to melt the gallium with hot water (m.p. $29.75^{\circ}$ ) and then to use a hypodermic needle to squirt the gallium through the $1 / 1 s^{\prime \prime}$ orifice into the cell. The vapor pressure of quartz was negligible at the highest temperature of these measurements.

Contrary to the observations of Harteck, no reaction of the gallium with quartz was noted. The surface of the gallium remained silvery and uncontaminated with the dark film observed by Harteck. The quartz liner was broken and examined to check this point at the conclusion of this research.
A disappearing filament optical pyrometer (Leeds and Northrup No. 8622) was used to measure the temperatures. The pyrometer was calibrated against a standard tungsten ribbon filament lamp standardized by the National Bureau of Standards. The precision of the temperature measurements is $0.1 \%$.

Sample. -The sample was weighed with an accuracy of 0.02 mg .

The gallium used was obtained from the Eagle-Picher Lead Company. On the basis of spectroscopic analysis, the gallium contained 0.001 part of lead and 0.002 part copper. This corresponds to a purity of 99.75 mole per cent.
It is possible to estimate the purity of the sample from premelting data obtained in determining the heat capacity in a low temperature calorimeter. ${ }^{4}$ A purity of 99.96 mole per cent. is obtained by this method. Before making any vapor pressure measurements, the gallium was heated at $1350^{\circ} \mathrm{K}$. for several hours in a high vacuum to rid the sample of volatile impurities.

## Data and Discussion

The vapor pressure can be calculated from effusion data using the formula

$$
\begin{equation*}
p=m / k \sqrt{2 \pi R T / M} \tag{1}
\end{equation*}
$$

[^0]where $p$ is the pressure in atmospheres, $m$ is the rate of evaporation in $\mathrm{g} . / \mathrm{sec} . / \mathrm{cm}^{2}{ }^{2}, T$ is the absolute temperature, $M$ is the molecular weight of the vapor, $R$ is the molar gas constant and $k$ is the correction due to the resistance to free molecular flow owing to the thickness of the orifice. For an orifice radius of 0.55 mm . and thickness 0.40 mm ., $k=0.736{ }^{5}$

The experimental results are recorded in Table I. These pressures are $59 \%$ higher than those recorded by Harteck, although Harteck's data scatter sufficiently so that some of his data overlap those obtained in this research. The reason for this difference may be a consequence of the difference in purity of the gallium used or may, perhaps, be due to the fact that the gallium amalgamated with the platinum and chrom-nickel Knudsen cells that Harteck employed in his research.

Table I
The Vapor Pressure of Gallium ${ }^{\text {a }}$

| $T,{ }^{\circ} \mathrm{K}$. | Mass <br> g. <br> $\times 10^{3}$ | Time, <br> sec. | $\mathrm{m}, / \mathrm{cm},{ }^{2} / \mathrm{sec}$. | Pressure, <br> atm. |
| ---: | ---: | ---: | ---: | :---: |
| 1230 | 1.73 | 14,752 | $1.675 \times 10^{-5}$ | $1.588 \times 10^{-6}$ |
| 1234 | 5.98 | 31,265 | $2.733 \times 10^{-5}$ | $2.593 \times 10^{-6}$ |
| 1255 | 2.74 | 13,924 | $2.811 \times 10^{-5}$ | $2.690 \times 10^{-6}$ |
| 1267 | 2.64 | 5,918 | $6.374 \times 10^{-5}$ | $6.130 \times 10^{-6}$ |
| 1274 | 7.17 | 12,863 | $7.964 \times 10^{-5}$ | $7.678 \times 10^{-6}$ |
| 1284 | 3.18 | 7,036 | $6.457 \times 10^{-5}$ | $6.251 \times 10^{-6}$ |
| 1304 | 4.15 | 4,848 | $1.223 \times 10^{-4}$ | $1.193 \times 10^{-5}$ |
| 1309 | 4.03 | 6,931 | $8.308 \times 10^{-5}$ | $8.119 \times 10^{-6}$ |
| 1317 | 4.00 | 4,903 | $1.166 \times 10^{-4}$ | $1.143 \times 10^{-5}$ |
| 1327 | 3.29 | 3,318 | $1.417 \times 10^{-4}$ | $1.395 \times 10^{-5}$ |
| 1330 | 1.70 | 1,493 | $1.626 \times 10^{-4}$ | $1.602 \times 10^{-5}$ |
| 1345 | 3.54 | 1,740 | $2.907 \times 10^{-4}$ | $2.880 \times 10^{-5}$ |
| 1370 | 2.90 | 1,250 | $3.313 \times 10^{-4}$ | $3.313 \times 10^{-5}$ |
| 1372 | 7.52 | 3,122 | $3.442 \times 10^{-4}$ | $3.444 \times 10^{-5}$ |
| 1380 | 5.45 | 1,395 | $5.581 \times 10^{-4}$ | $5.601 \times 10^{-5}$ |
| 1385 | 9.71 | 3,367 | $4.120 \times 10^{-4}$ | $4.142 \times 10^{-5}$ |
| 1405 | 18.79 | 3,515 | $7.638 \times 10^{-4}$ | $7.733 \times 10^{-5}$ |
| 1425 | 21.51 | 3,369 | $9.122 \times 10^{-4}$ | $9.301 \times 10^{-5}$ |
| 1518 | 18.52 | 523 | $4.122 \times 10^{-3}$ | $4.339 \times 10^{-4}$ |

${ }^{a}$ These vapor pressures have been computed upon the assumption that gallium vapor is monatomic.

The reliability of the vapor pressure data is best tested by establishing the invariance of $\Delta H_{0}^{8}$ (the heat of sublimation at absolute zero) in the thermodynamic relationship


[^1]

Fig. 1,--Gallium: $\log P$ as. $1 / I$.
when applied to the individual measurements of the equilibrium

$$
\begin{equation*}
\text { Ga (liquid) }=\text { Ga (vapor) } \tag{B}
\end{equation*}
$$

The thermodynamic functions for the liquid are given by the expression

$$
\begin{equation*}
\left(\frac{F^{0}-H_{0}^{0}}{T}\right)_{1 \mathrm{i}, \mathrm{p}}=\frac{1}{T} \int_{0}^{T} C_{\mathrm{p}} \mathrm{~d} T-\int_{0}^{T} C_{\mathrm{y}} \mathrm{~d} \ln T \tag{4}
\end{equation*}
$$

The following thermal data were used to calculate the thermodynamic functions for liquid gallium listed in column 2 of Table II.

Table II
Thermodynamic Calculations for Gallium

| ( ${ }^{\mathrm{FO}}$ | $\begin{aligned} & )_{\substack{\text { img } \\ \text { cal.g. } \\ \text { atom/des. }}} .- \end{aligned}$ |  | $\begin{gathered} -k \ln \\ p(\mathrm{atm.m}), \\ \text { cal./g. } \\ \mathrm{atom} / \mathrm{deg} . \end{gathered}$ | $\begin{gathered} \text { Alli, } \\ \text { gual } \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: |
| 1239 | 16.4; | +4.04 | 26.53 | 66.6 |
| 12.44 | 16.5 | 4.05 | 26.56 | 66.8 |
| 1255 | 16.57 | 4.15 | 25.49 | 66.6 |
| 1267 | 16.64 | $4 \pm .19$ | 23.85 | 03.1 |
| 1274 | 16.68 | 44.22 | 23.41 | 64.9 |
| 1284 | 16.74 | 44.26 | 23.82 | 63.9 |
| 1304 | 16.85 | 44.35 | 22.3 | (6). 2 |
| 1309 | 16.87 | 44.5 | 23.30 | (6). 5 |
| 1317 | 16.92 | 44.40 | 22.61 | (6). 0 |
| 1327 | 16.98 | 44.4 | 22.10 | 68.9 |
| 1330 | 16.99 | 4.45 | 21.9 - | (5). 7 |
| 1345 | 17.07 | 4.51 | 20.78 | 040 |
| 1370 | 17.21 | $4+60$ | 20. 01 | (65. i |
| 1372 | 17.22 | 4.6.61 | 20.42 | 6 6 |
| 1380 | 17.26 | 4.64 | 19.46 | 64.6 |
| 1385 | 17.28 | 4t.6it | 20.66 | (35. $\%$ |
| 1405 | 17.38 | 44.74 | 18.42 | 64.4 |
| 1425 | 17.49 | +1.8.81 | 18.45 | (0.3.2 |
| 1518 | 17.94 | 45.16 | 15.4\% | 64.7 |
|  |  |  | Average | (ib $t=0$ |

Values for the thermodynamic functions up to $320^{\circ} \mathrm{K}$. were taken from previous calculations made in this Laboratory. These functions were extended into the high temperature range using a value of $6.645 \mathrm{cal} . / \mathrm{g}$. atom $/ \mathrm{deg}$. for the heat capacity of liquid gallium (determined at $320^{\circ} \mathrm{K}$.). The heat
capacity of liquid gallium was assumed to be constant over the entire liquid range, and as a consequence there may be some error in the thermal functions for the liquid; however, this error should be quite small because the heat capacity for many liquid metals is virtually constant.

The free energy function for gallium vapor depends both on the vapor pressure and the energy states of the gallium vapor. The free energy function for the vapor is

$$
\begin{array}{r}
\left(\frac{F^{\nu}-i!\prime}{T}\right)_{\text {anor }}=-3 / 2 R \ln M-5 / 2 R \ln T-C- \\
R \ln R-R \ln Q(5) \tag{5}
\end{array}
$$

where $M$ is the molecular weight of the vapor, $C$ $+R \ln R=-7.267$ calories if $p$ is expressed in atmospheres, $Q=\sum_{i} g_{i} e^{-\varepsilon_{i} / k T}$ is the partition function and $g_{i}$ is the a priori weight of the state of energy $\epsilon_{i} ; k$ is Boltzmann's constant.

All terms but the last on the right-hand side of equation (5) are the translational contribution to the free energy function while the last term is the contribution due to the distribution of the gallium atoms among the available energy states of gallium at a given temperature. In the temperature range of these experiments only the ground state, ${ }^{2} \mathrm{P}^{0}{ }_{1 / 2}$, and the first excited state, ${ }^{2} \mathrm{P}^{0^{0} / 2}$, make an appreciable contribution to the internal energy of the vapor. Consequently

$$
\begin{equation*}
Q=2+4 \mathrm{e}^{-\epsilon_{2} / k T} \tag{6}
\end{equation*}
$$

where $\epsilon_{1}=0$ is the energy of the ground state and $\epsilon_{2}=\$ 26.6 \mathrm{~cm} .^{-1}$, is the energy of the first excited state.

The free energy functions for gallium vapor are given in column 3 of Table II.

The $\Delta H_{0}^{0}$ values computed by equation (2) are given in the last column of Table II. These values show no temperature trend and are constant to within the experimental error of the vapor pressure data. The average value of $65.6 \pm 0.5$ kcal. for $\Delta H 0$ is about 2 kcal. lower than the value 67.7 calculated from the data of Harteck. ${ }^{2}$ The heat of vaporization at the mean temperature $\left(13.50^{\circ} \mathrm{K}\right.$.) of these experiments is

$$
\begin{aligned}
\Delta H_{1,3}^{1},\left(H_{1300}^{1}\right)_{g}-\left(H_{13,0}^{0} h_{1}+\Delta H_{0}^{0}\right. & =6.7-9.7+ \\
65.6 & =62.6 \mathrm{kcal} .
\end{aligned}
$$

Least squares treatment of our own data yields the following vapor pressure equation for gallium
$\operatorname{lug} p(a t m)=.8.540-\frac{1.200}{T} \times 10^{2}-0.844 \log T$

The boiling point of gallium, as computed from equation (8), is $2500^{\circ} \mathrm{K}$. or $2227^{\circ} \mathrm{C}$.

This equation yields a value of 60.7 kcal . for the heat of vaporization of liquid gallium which differs by 1.9 kcal . from the value obtained in equation (7).

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[^0]:    (1) This work was carried out under contract between the Office of Naval Research and The Ohio State University Research Foundation.
    (2) P. Harteck, Z. physik. Chem., 134, 1 (1928).
    (3) R. B. Holden, R. Speiser and H. L. Johnston, This Journal, 70, 3897 (1948); R. Speiser and H. L. Johnston. Preprint A.S.M., Cleveland, Ohio, Oct. 17, 1949.
    (4) G. B. Adams, H. L. Johnston and E. C. Kerr, This Journal, 74, 4784 (1952).

[^1]:    (5) S. Dushman, 'Scientific Foundations of Vacuum Technique,'" John Wiley and Sons, Inc., New York, N. Y., 1949, p. 99.

